

# The search and detection of quantum spin liquid in new materials with geometrically frustrated lattice

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*09-19-2022*

# Acknowledgements:

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- **University of Tennessee**  
Jian Liu, Hang Zhan, Qing Huang, XingChen Kun
- **University of Science and Technology, China**  
XueFeng Sun
- **National High Magnetic Field Lab**  
Eunsang Choi
- **University of Michigan**  
Lu Li



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**ENERGY**

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Science

# Motivation:

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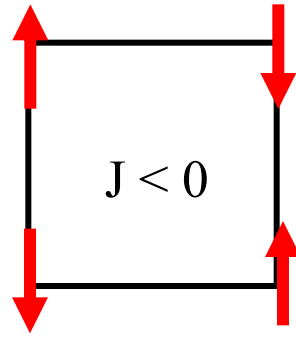
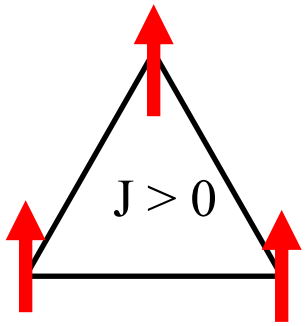
*We must learn how the astonishing properties of quantum materials can be tailored to address our most pressing technological needs, and we must dramatically improve our ability to synthesize, characterize, and control quantum materials.*

*Grow single crystals*

*Search for quantum spin liquid state candidates (NSF)*

*Electronically detect spin states and magnetic excitations (DOE, collaboration with Jian Liu)*

# Non-frustrated magnets



$$H = - \sum_{ij} J_{ij} S_i \cdot S_j$$

$$T_N = \frac{zS(S+1)|J|}{3k_B}$$

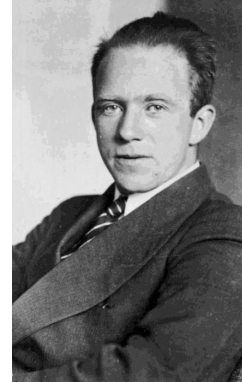
$$\chi = \frac{C}{T - \theta_{CW}}$$



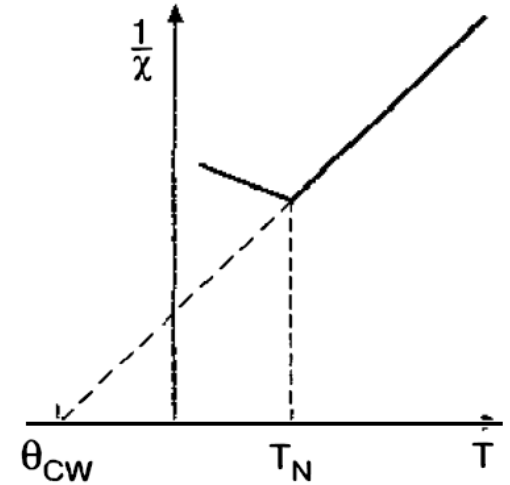
Curie



Néel



Heisenberg



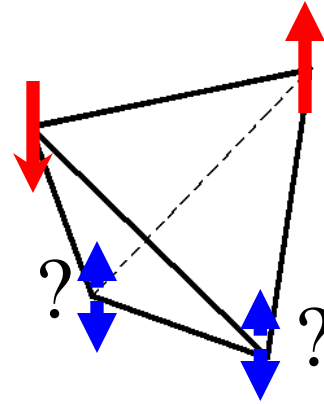
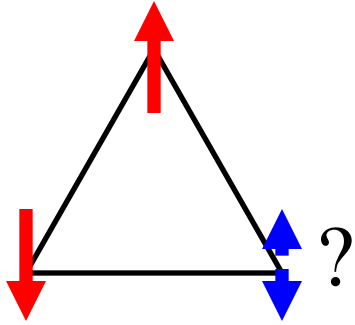
Non-frustrated  $T_N \sim \theta_{CW}$

$z = 4$  nearest neighbor number

$S = 1$

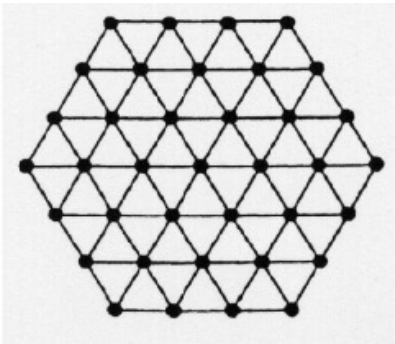
$T_N \sim 2.7 J$

# Geometrically Frustrated Lattice

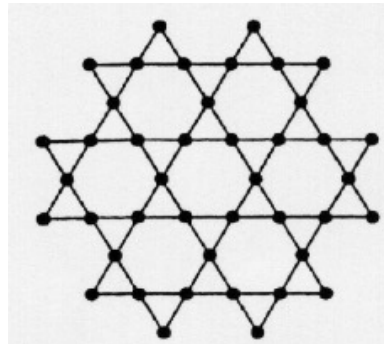


*Interactions between magnetic degree of freedom in a lattice are incompatible with the underlying crystal geometry -----Frustration*

2D

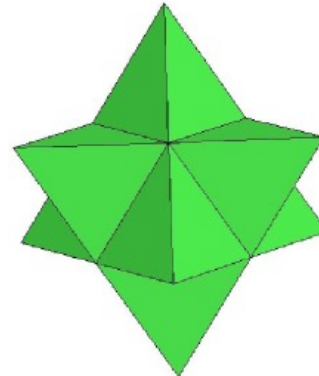


Triangular

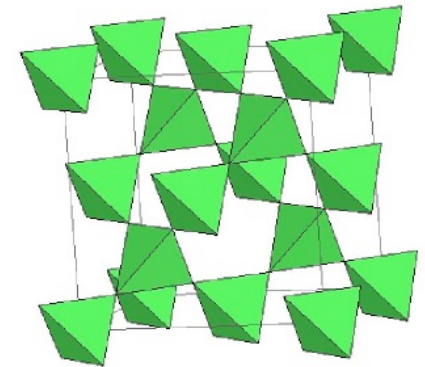


Kagome

3D

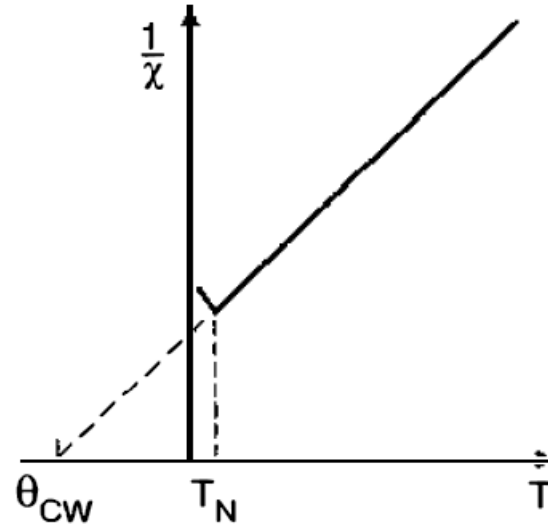
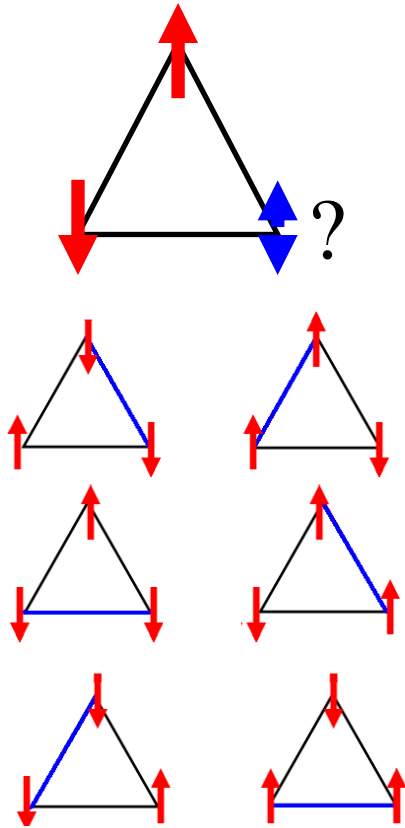


Face centered cubic



Pyrochlore

# Degeneracy

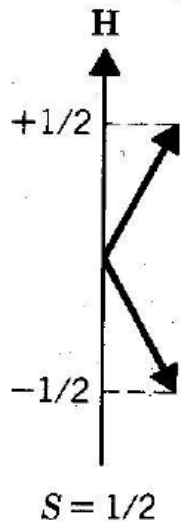


Frustrated  $T_N \ll \theta_{CW}$

$$f = |\theta_{CW}|/T_N > 10$$

Frustration leads to degeneracy, which enhances spin fluctuations and suppresses magnetic ordering to *induce exotic magnetism*.

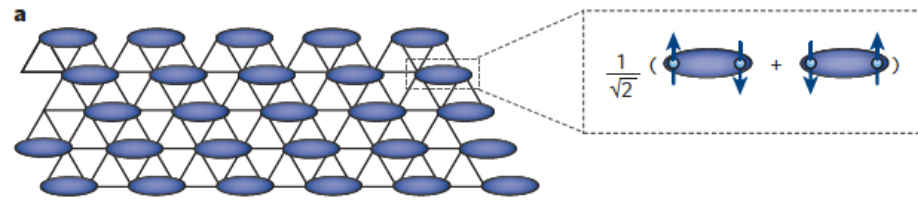
# Spin-1/2 Triangular lattice antiferromagnet (TLAF)



*Strong quantum spin fluctuations*

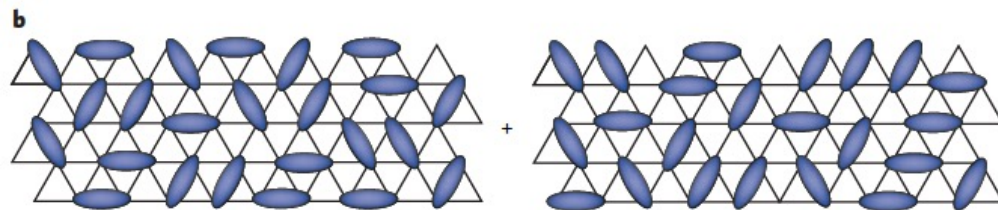
# Quantum spin liquid

**Quantum spin liquid (QSL) RARE!!** No long range order down to 0 K;  
 No symmetry breaking; Long range entanglement; Fractional excitation



$$(i, j) = \frac{1}{\sqrt{2}} (|\uparrow_i \downarrow_j\rangle - |\downarrow_i \uparrow_j\rangle)$$

Valence-bond solid (VBS) state: a singlet dimer configuration dominates in the ground state.



$$|\Psi_{\text{RVB}}\rangle = \sum_{i_1 j_1 \dots i_n j_n} a_{(i_1 j_1 \dots i_n j_n)} |(i_1, j_1) \dots (i_n, j_n)\rangle,$$

Resonant valence bond (RVB) state: The valence-bond pairs in the RVB construction are dominated by short-range pairs, resulting in liquid-like states with no long-range spin order.

*P. W. Anderson, Mater. Res. Bull. 8, 153 (1973); Science 235, 1196 (1987)*

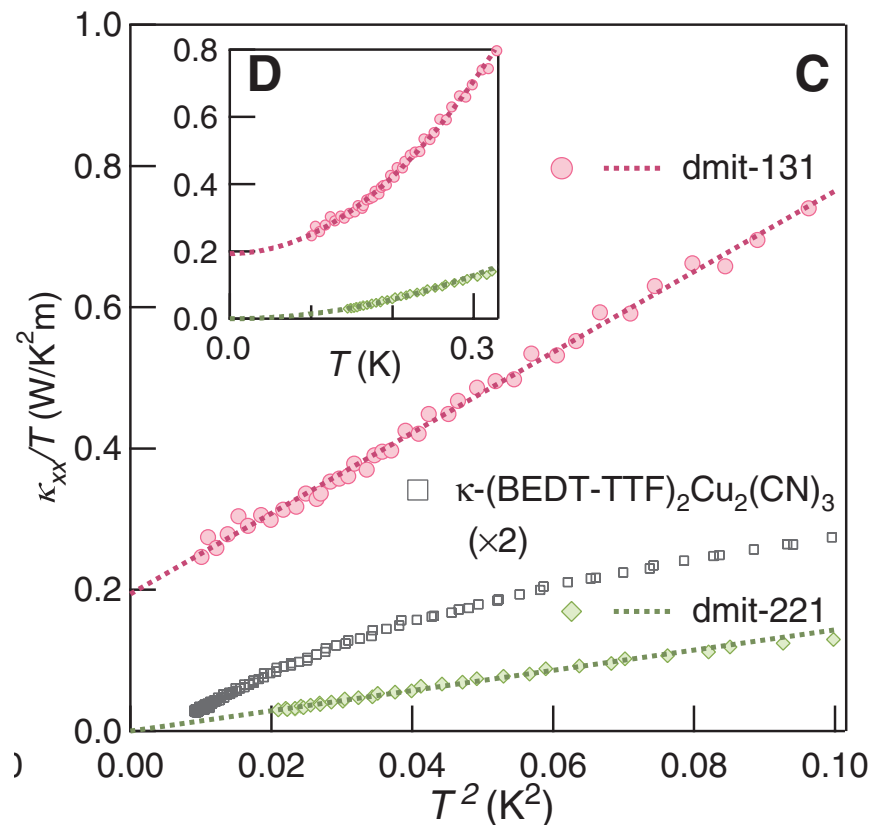
*L. Balents, Nature 464, 199 (2010)*



# Gapless QSL, itinerant spin excitations, residual $k_0/T$ term

Organic molecular magnets:  $\text{EtMe}_3\text{Sb}[\text{Pd}(\text{dmit})_2]_2$ , spin-1/2 TLAF

*M. Yamashita et al., Science 4, 328 (2010).*



Very rare to observe non-zero  $k_0/T$  approaching zero temperature in insulating magnets

# Spin-1/2 geometrically frustrated magnets

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$\text{Cu}^{2+}, 3d^9$

$\text{ZnCu}_3(\text{OH})_6\text{Cl}_2$ , Herbertsmithite, kagome lattice, QSL

$\text{Cs}_2\text{CuBr}_4$ , distorted triangular lattice, LRO with UUD

$\text{Ir}^{4+}, 5d^5$

$\text{Na}_4\text{Ir}_3\text{O}_8$ , hyper-kagome, QSL

$\text{Ru}_2\text{O}_9$  dimer

$\text{Ba}_3\text{ARu}_2\text{O}_9$  (A =  $\text{Y}^{3+}$ ,  $\text{In}^{3+}$ ,  $\text{Lu}^{3+}$ )

$\text{Co}^{2+}, 3d^7$ , effective spin 1/2

$\text{Ba}_3\text{CoSb}_2\text{O}_9$ ,

$\text{B}_3\text{CoNb}_2\text{O}_9$ ,

$\text{Ba}_8\text{CoNb}_6\text{O}_{24}$

$\text{Mo}_3\text{O}_{13}$  clusters

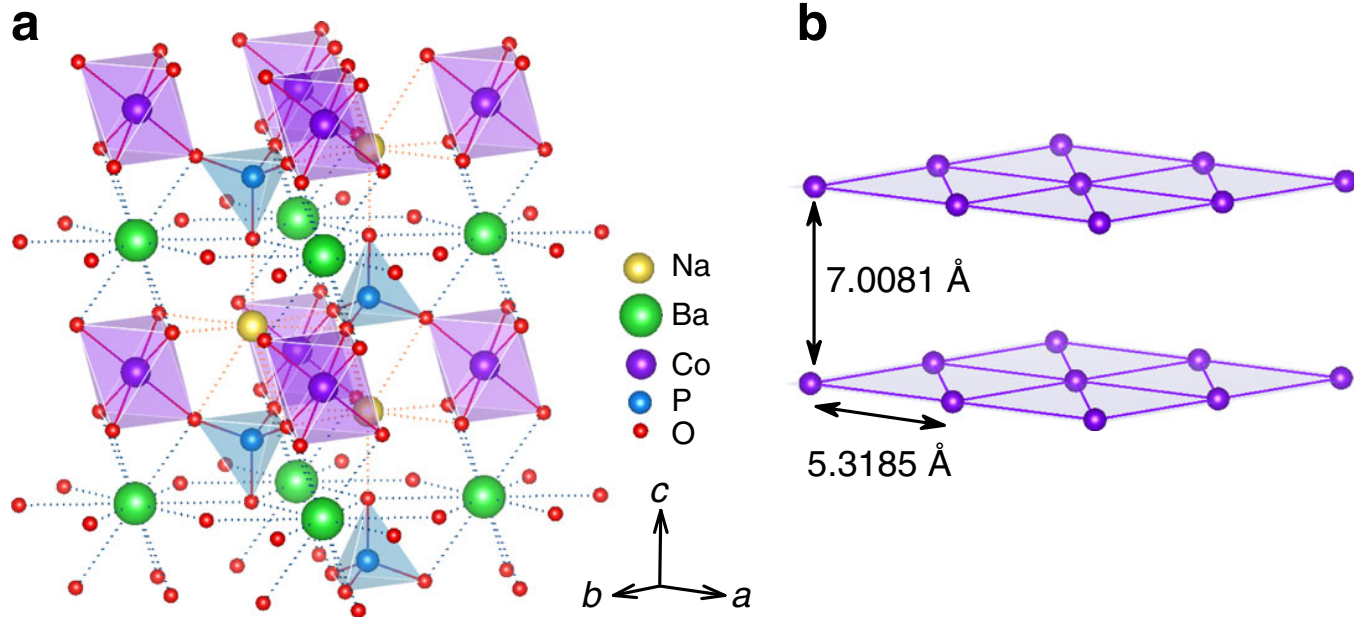
$\text{LiZn}_2\text{Mo}_3\text{O}_8$ , distorted Kagome lattice QSL

$\text{Li}_2\text{In}_{1-x}\text{Sc}_x\text{Mo}_3\text{O}_8$

$\text{Yb}^{3+}, 4f^{13}$ , effective spin 1/2

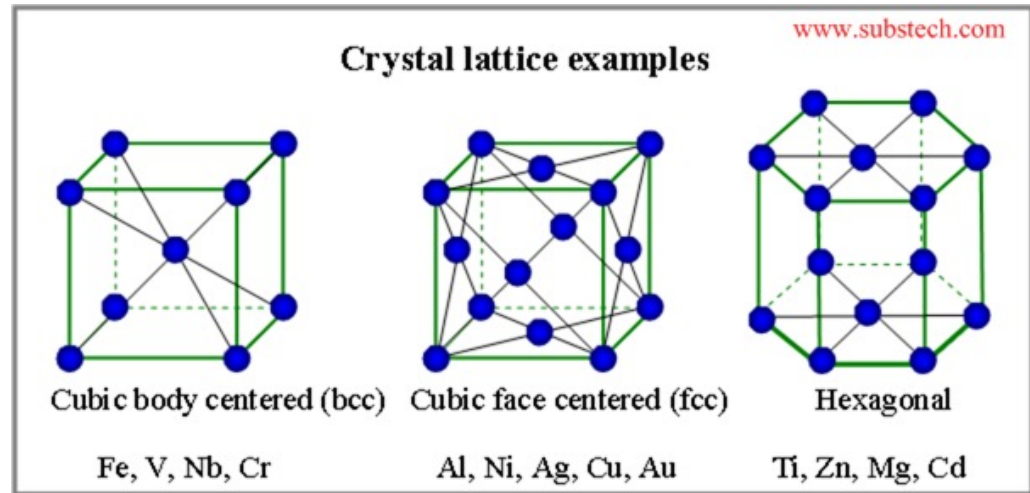
$\text{Yb}_2\text{Ti}_2\text{O}_7$ , pyrochlore, QSL

# $\text{Na}_2\text{BaCo}(\text{PO}_4)_2$ , spin-1/2 triangular lattice antiferromagnet



# Single Crystal

A **single crystal** is a material in which the crystal structure of the entire sample is continuous and unbroken to the edges of the sample, with no grain boundaries.



# Crystallization

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*Nucleation* is the step where the solute molecules dispersed in the solvent start to gather into clusters, on the nanometer scale (elevating solute concentration in a small region)

The *crystal growth* is the subsequent growth of the nuclei that succeed in achieving the critical cluster size.



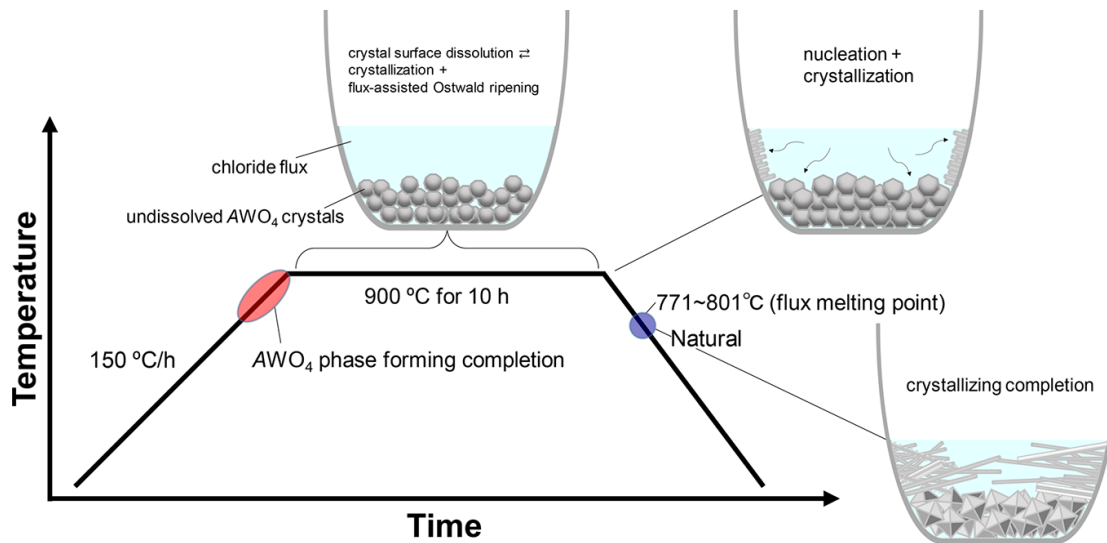
# Flux growth

e.g. Rock candy

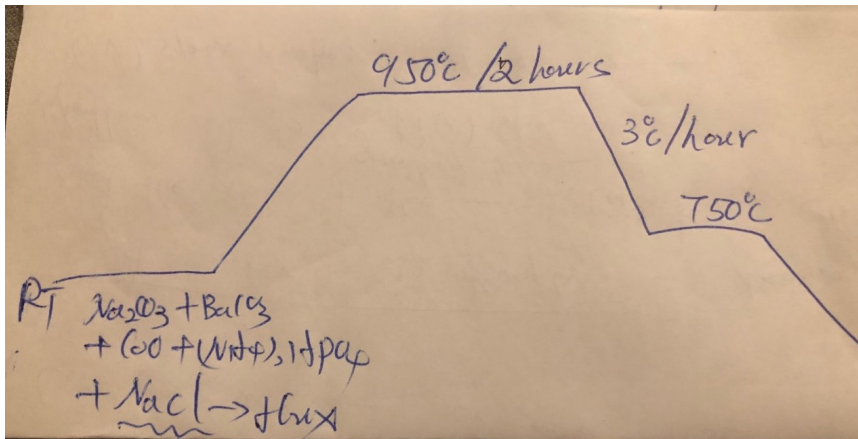


• Find the right solvent and dissolve the starting materials

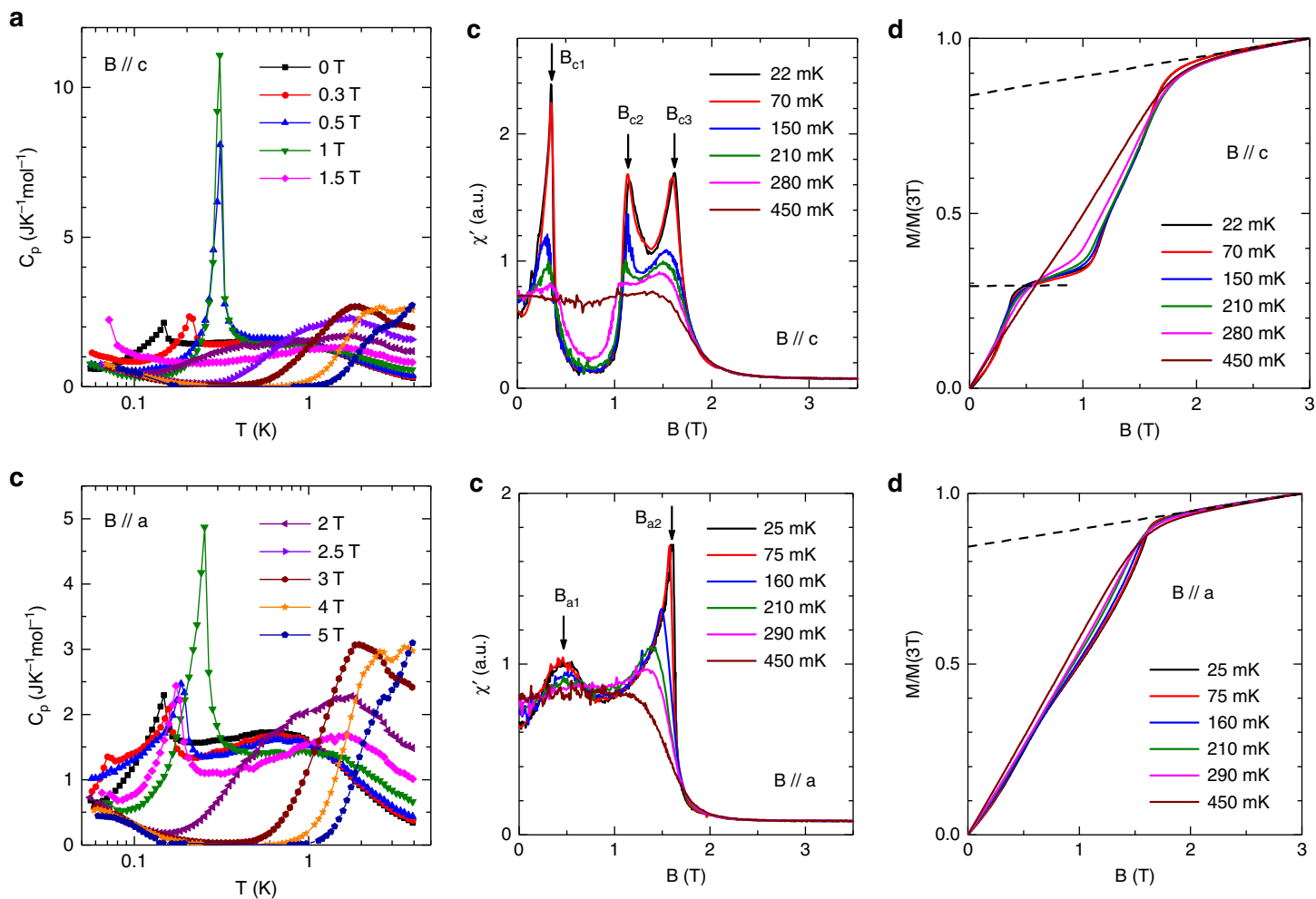
• Crystallize with time and temperature



# Flux growth



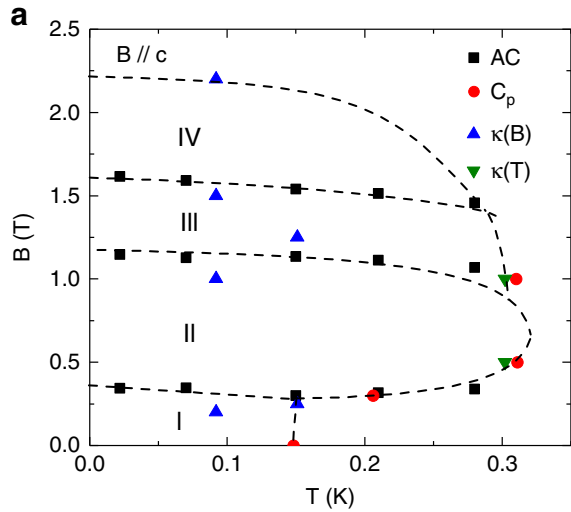
# $\text{Na}_2\text{BaCo}(\text{PO}_4)_2$ , $T_N = 0.15$ K, field induced spin state transitions



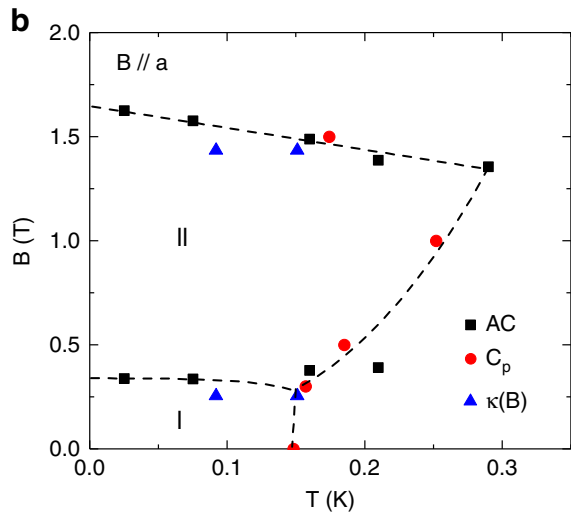
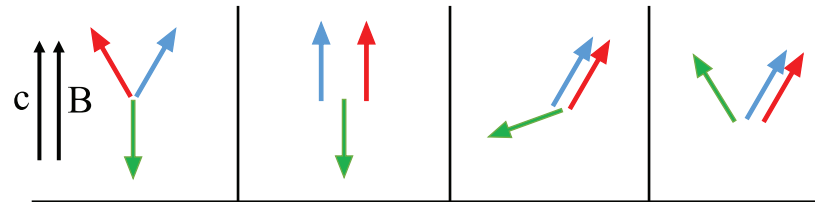
*N. Li, H. D. Zhou et al., Nature Communications 11, 4216 (2020).*



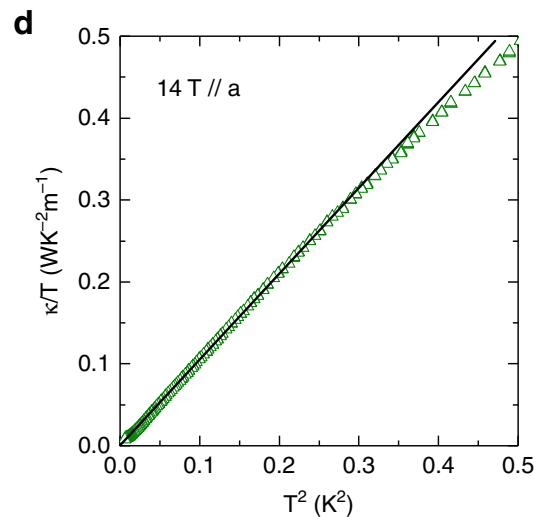
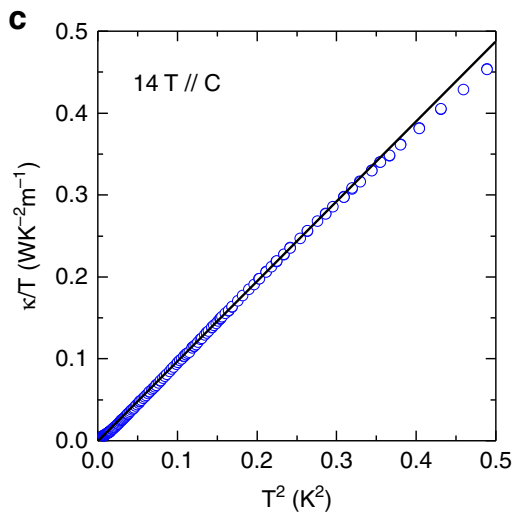
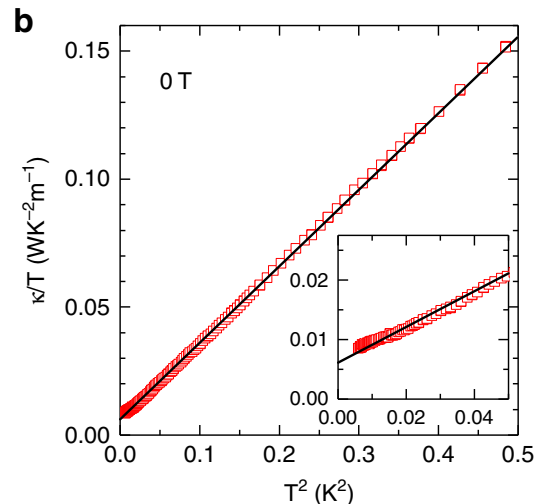
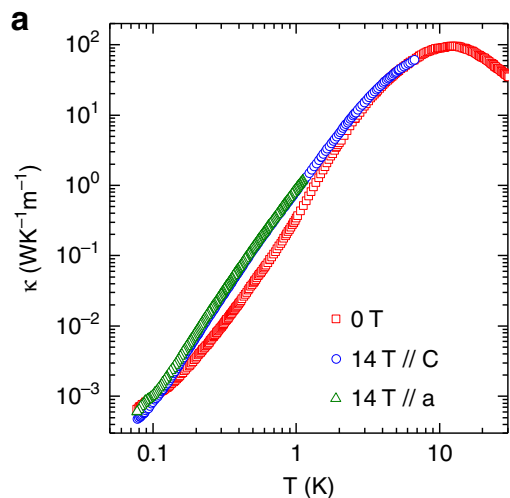
# $\text{Na}_2\text{BaCo}(\text{PO}_4)_2$ , magnetic phase diagram



the UUD phase only survives for  $B // c$ , which strongly suggests its easy axis anisotropy.



# $\text{Na}_2\text{BaCo}(\text{PO}_4)_2$ behaves as a gapless QSL above $T_N$

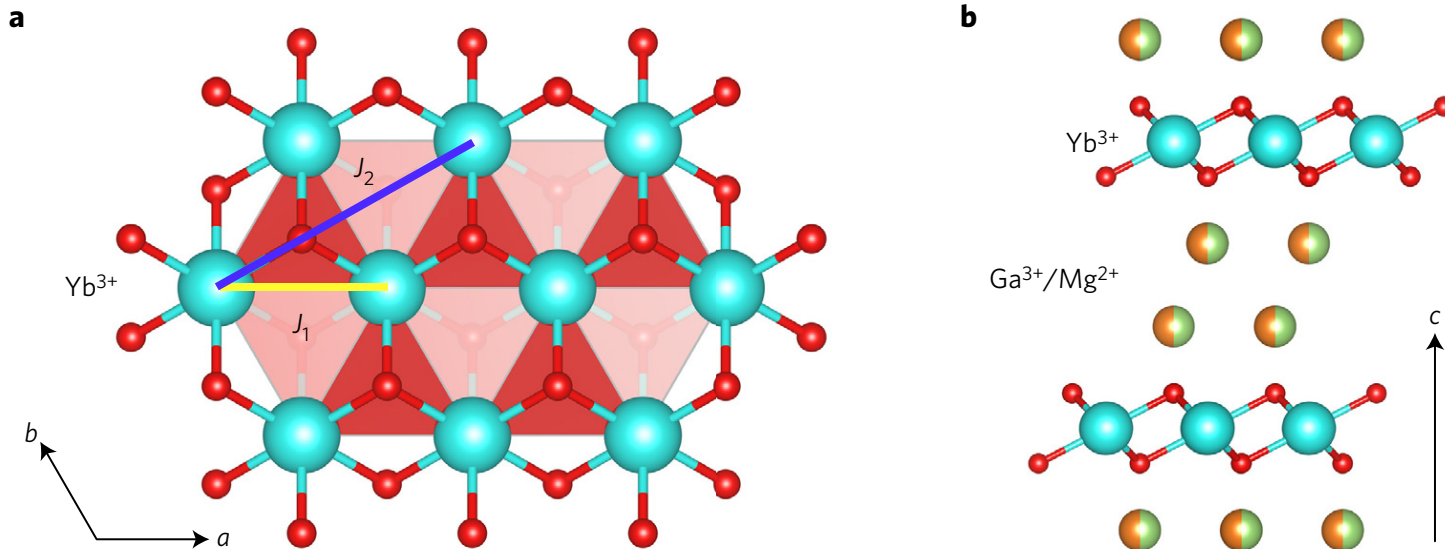


$$\frac{\kappa_0}{T} = \frac{\pi k_B^2 l_s}{9\hbar ad} = \frac{\pi}{9} \left(\frac{k_B}{\hbar}\right)^2 \frac{J}{d} \tau_s$$

$a \sim 5.32 \text{ \AA}$  and  $d \sim 7.01 \text{ \AA}$  are nearest-neighbor and interlayer spin distance, respectively.

From the observed  $\kappa_0/T = 0.0062 \text{ WK}^{-2}\text{m}^{-1}$ , the  $l_s$  (mean free path) is obtained as  $36.6 \text{ \AA}$ , indicating that the excitations (spinons) are mobile to a distance seven times as long as the inter-spin distance without being scattered.

# $\text{YbMgGaO}_4$ , spin-1/2 triangular lattice antiferromagnet

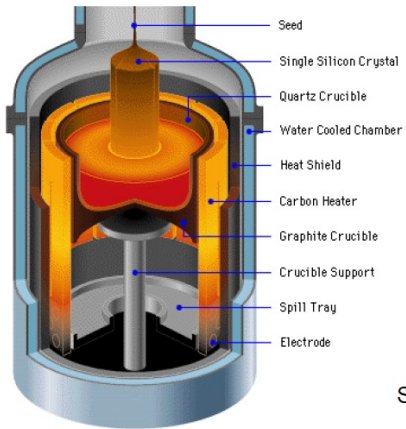


# Melting Growth

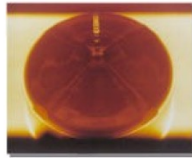


Bi Crystals grown in kitchen

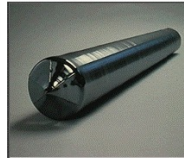
## Czochralski Crystal Growth



Crystal Pulling



Crystal Ingots



Shaping and Polishing



300 mm wafer

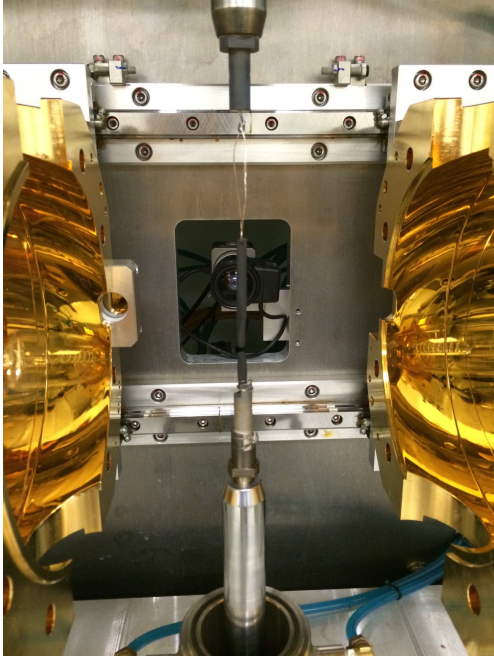


Si crystal growth

# Image Furnaces

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## Two mirrors Image Furnace

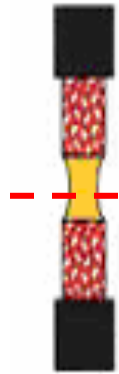


2200 Celsius degree

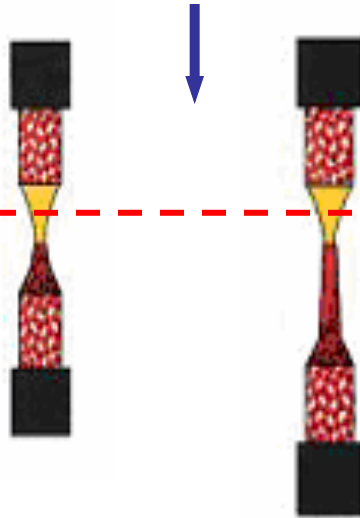


# Crystal growth, floating zone technique

stabilizing

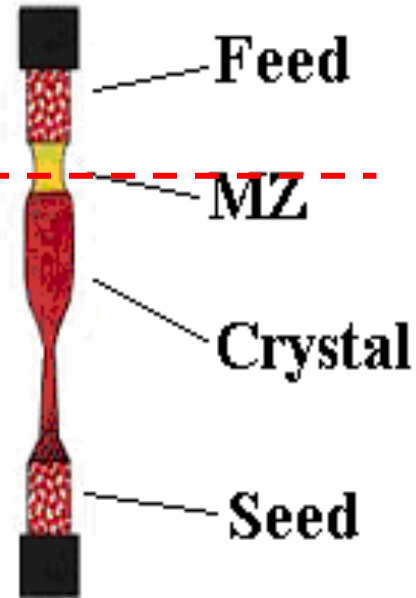


Necking



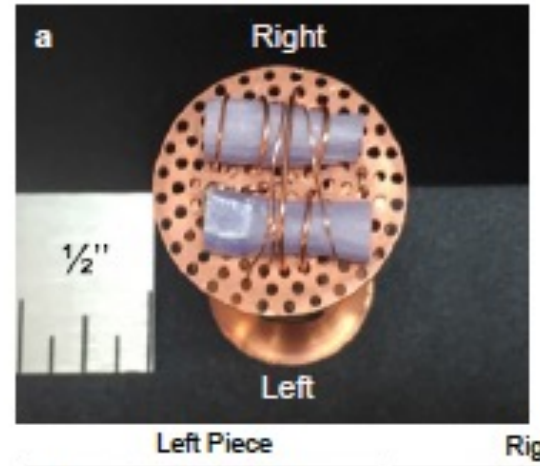
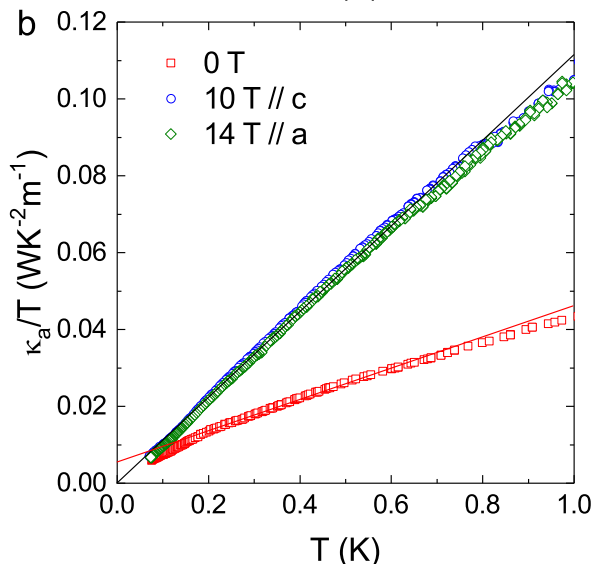
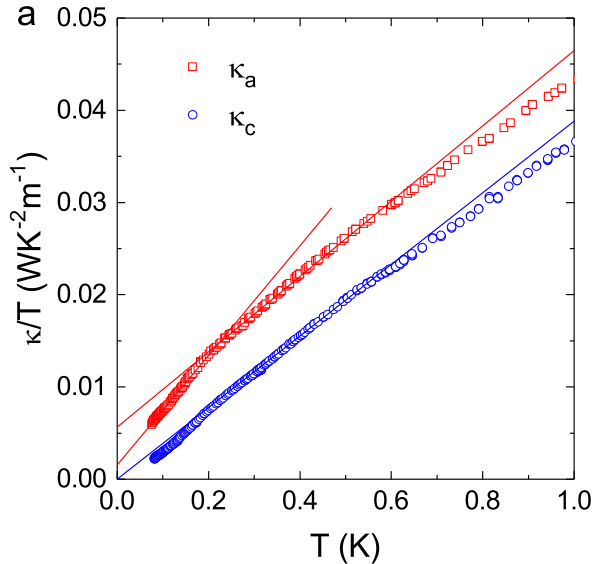
Feeding: 1 mm/h  
Growth: 3mm/h

Growth



Feeding: 2 mm/h  
Growth: 3mm/h

# YbMgGaO<sub>4</sub>, residual $\kappa_0/T$ term

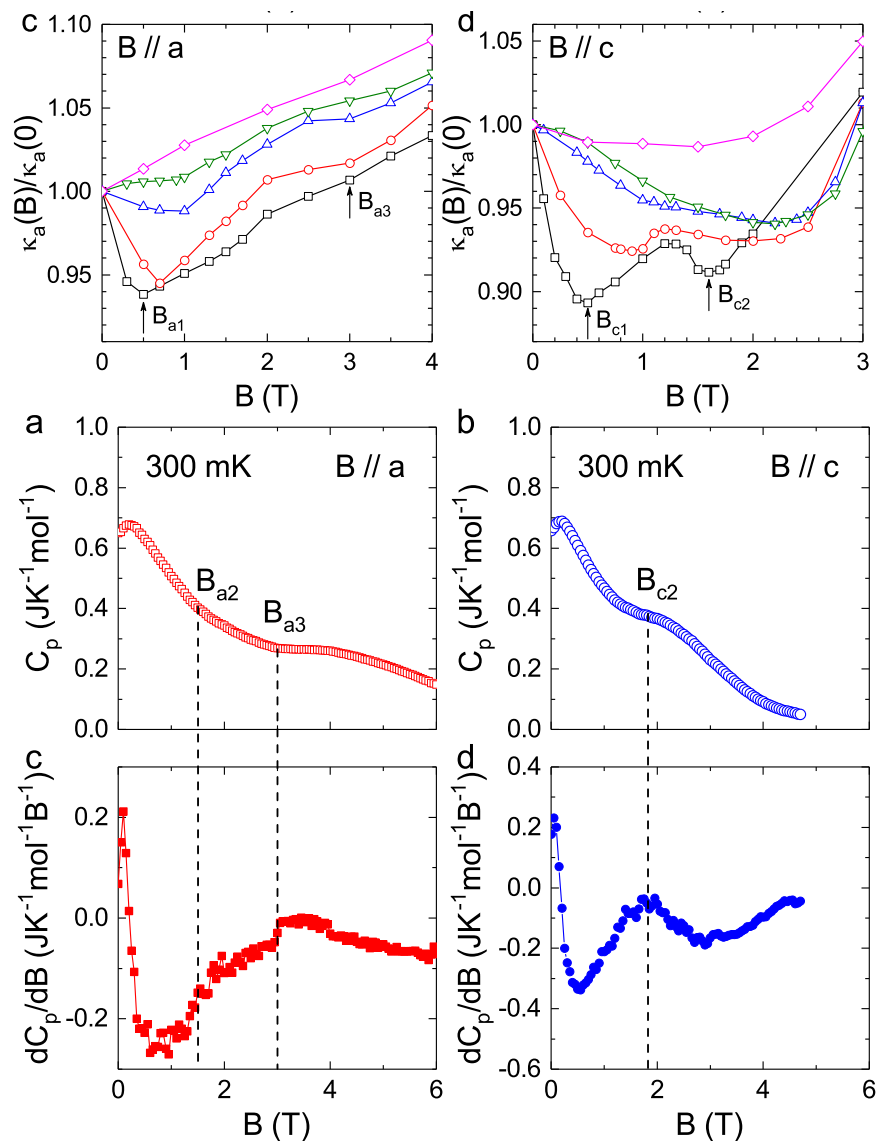


From the observed  $\kappa_0/T = 0.0058 \text{ WK}^{-2}\text{m}^{-1}$ , the  $l_s$  (mean free path) is obtained as  $78.4 \text{ \AA}$ , indicating that the excitations (spinons) are mobile to a distance 23 times as long as the inter-spin distance without being scattered.

The gapless QSL with itinerant excitations survives with Mg/Ga disorder

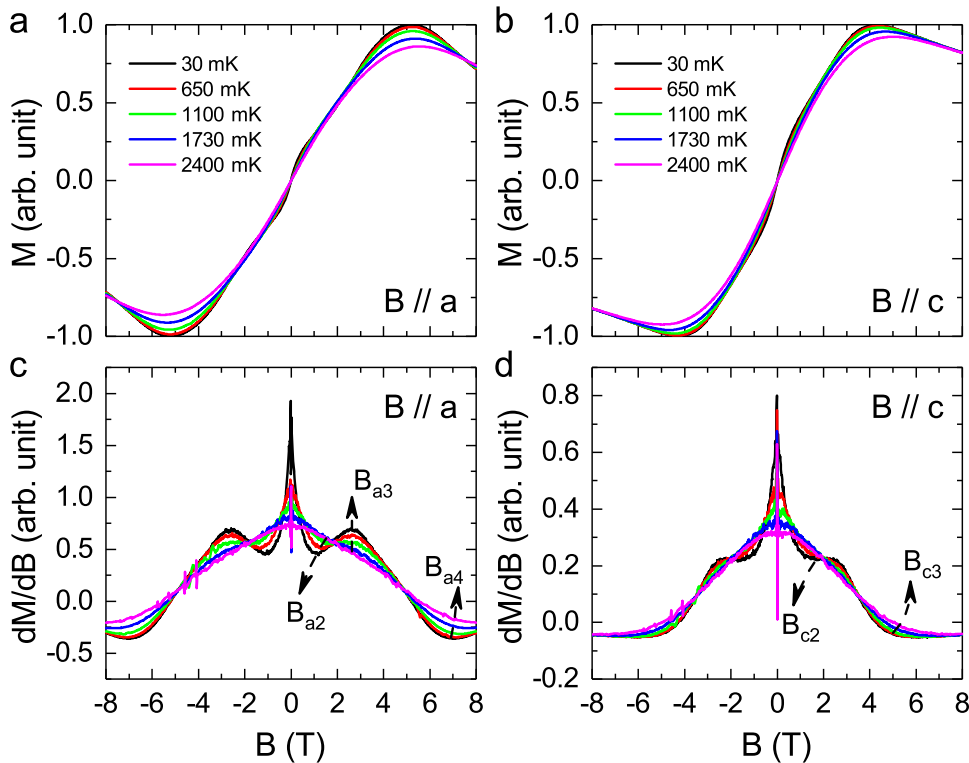
*X. Rao, H. D. Zhou et al., Nature Communications 12, 4949 (2021).*

# YbMgGaO<sub>4</sub>, field induced spin state transitions



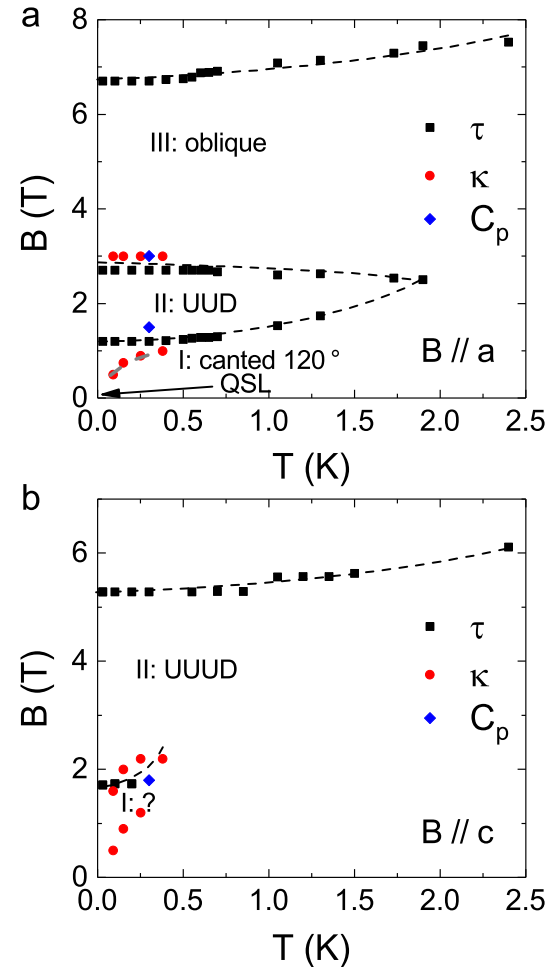


# YbMgGaO<sub>4</sub>, field induced spin state transitions



UUD phase with  $B // a$   
 UUUD phase with  $B // c$

Magnetization plateau feature is weak due to disorder



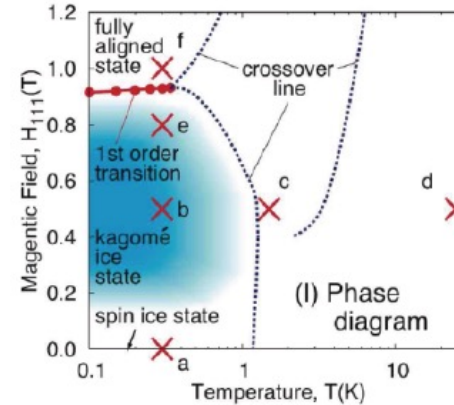
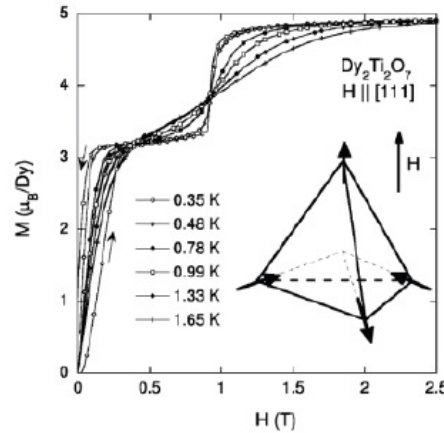
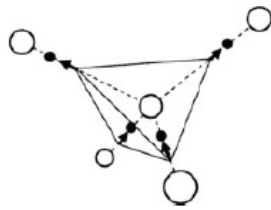
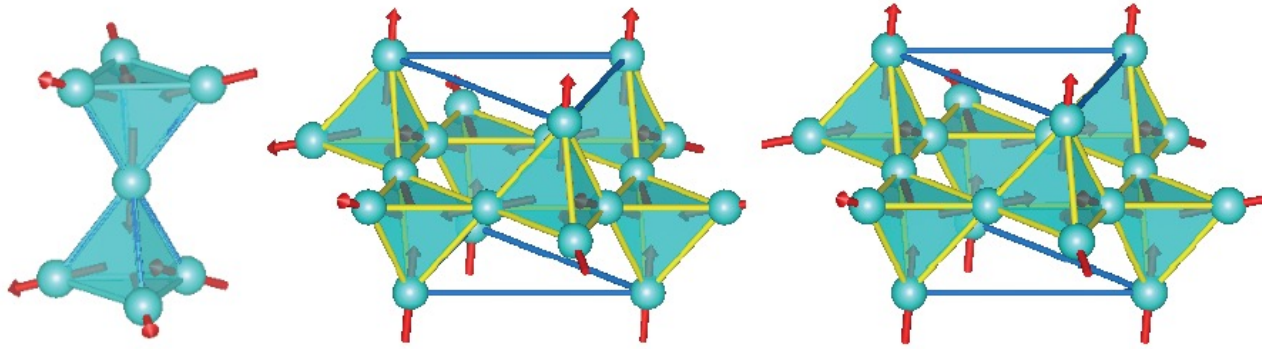
# Pyrochlore heterostructure

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*To explore a new route towards “metallization of quantum magnet” or electronically detect spin states and magnetic excitations*

*The hypothesis of the proposed approach is that, when combined to form a heterostructure, the interfacial coupling between the magnetic degree of freedom in an insulating geometrically frustrated quantum magnet (GFQM) and the electronic degree of freedom in a spin-orbit-entangled correlated metal necessarily leads to electronic transport signatures that are characteristic of the unusual spin states and their elementary excitations*

# Select Pyrochlore $\text{Dy}_2\text{Ti}_2\text{O}_7$ (DTO) as the GFQM

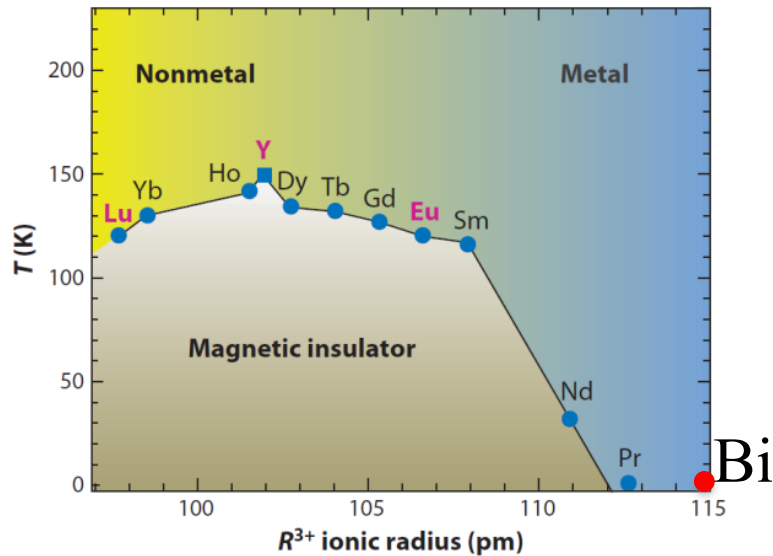


With applied field along the  $[111]$  axis, the spin ice (two in two out) state transforms to kagome spin ice and then three in one out state

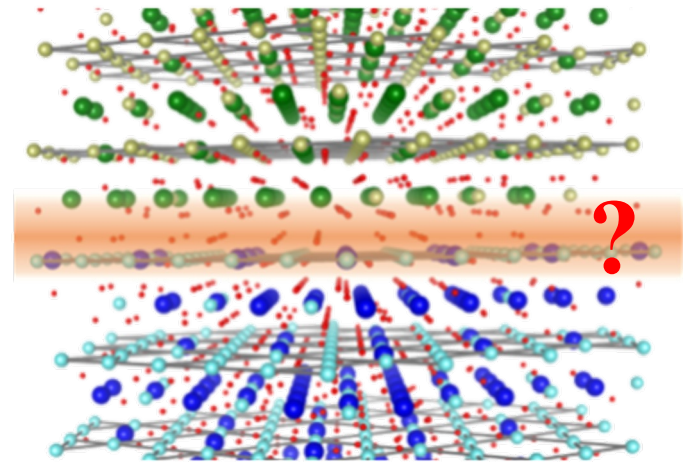
*Sakakibara, T. et al Phys. Rev. Lett. 90, 207205 (2003).*

*Tabata, Y et al., Phys. Rev. Lett. 97, 257205 (2006).*

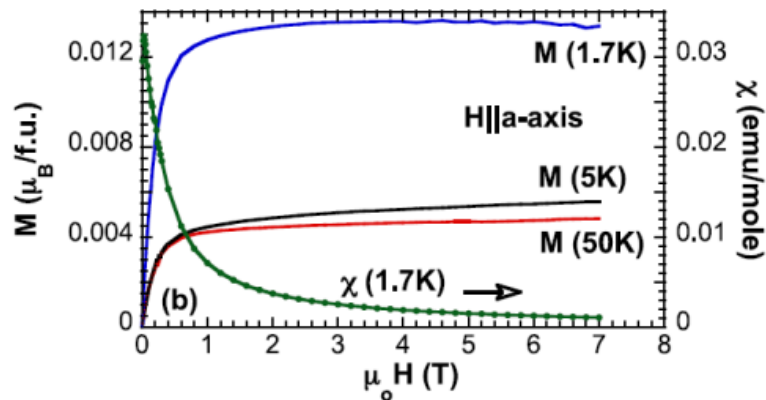
# Select Pyrochlore $\text{Bi}_2\text{Ir}_2\text{O}_7$ (BIO) as the Correlated Metal



Paramagnetic metal BIO



Spin Ice DTO

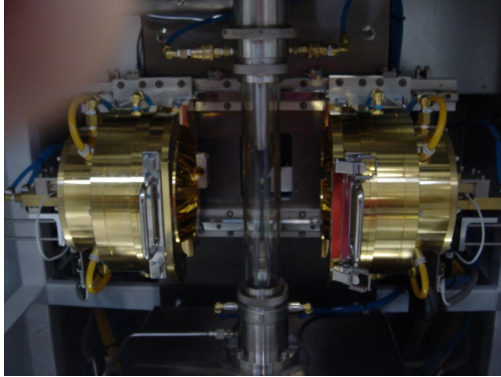


*T. F. Qi et al., J. Phys. Condens. Matter 24, 345601 (2012);*

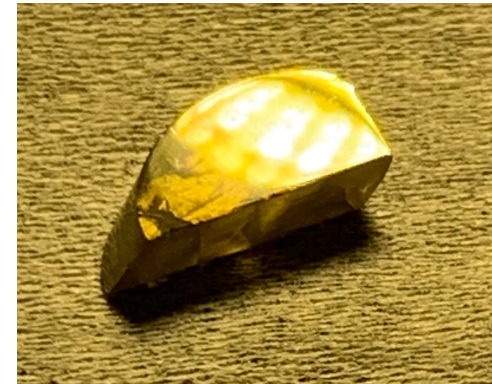
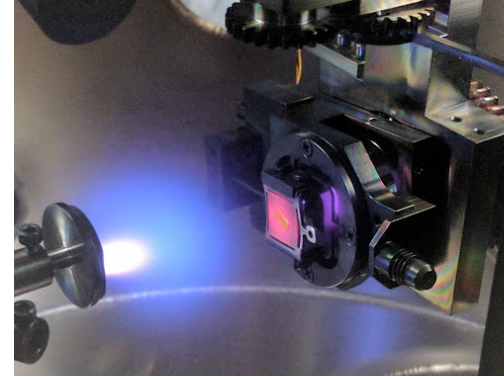
*W. Witczak-Krempa et al., Annual Review of Condensed Matter Physics 5, 57 (2014).*

# Growth (I): BIO film on DTO single crystal

## Floating Zone

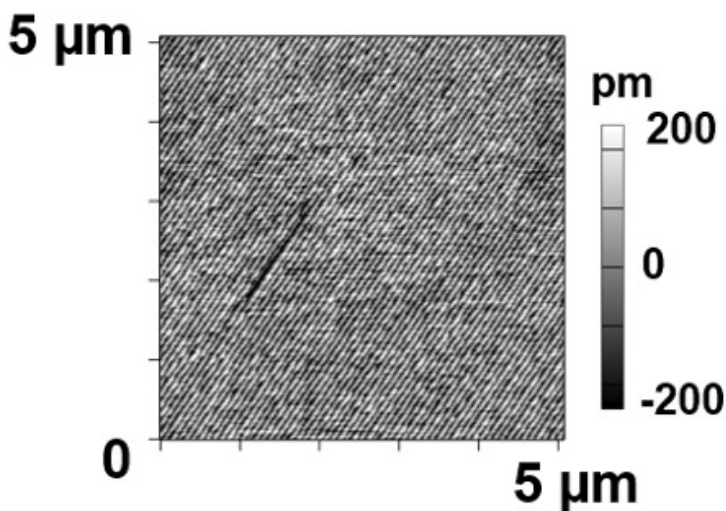
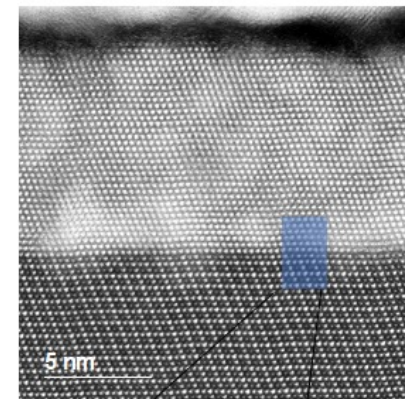
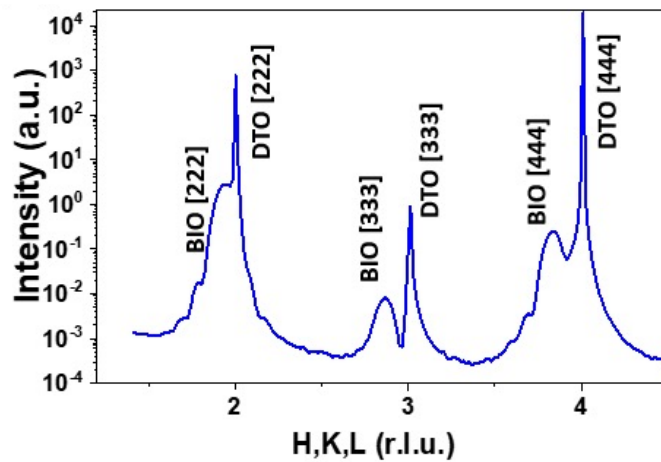
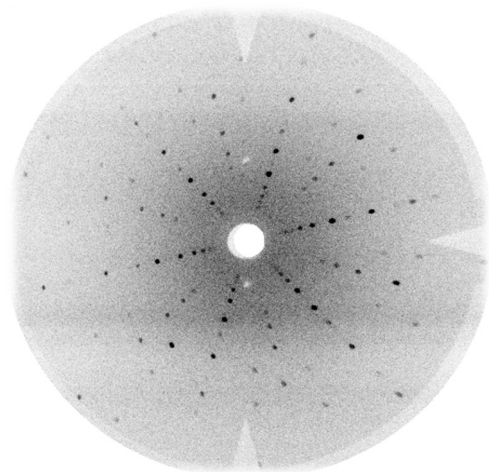


## Pulsed Laser Deposition

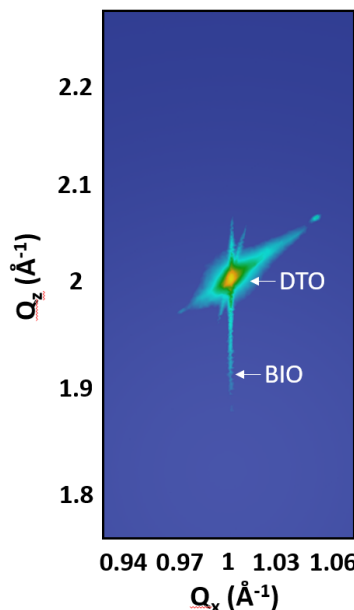


A combination of optical floating zone crystal growth and pulsed laser deposition is used to synthesize the DTO/BIO heterostructures along the [111] direction.

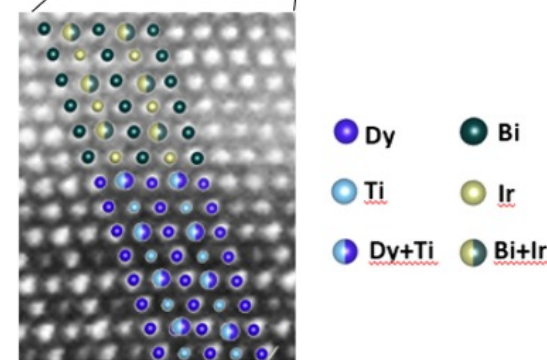
# Growth (II): Orientation, Polish, Film, Interface



AFM,  $R_s \sim 1.21 \text{ \AA}$



RSM, fully strained

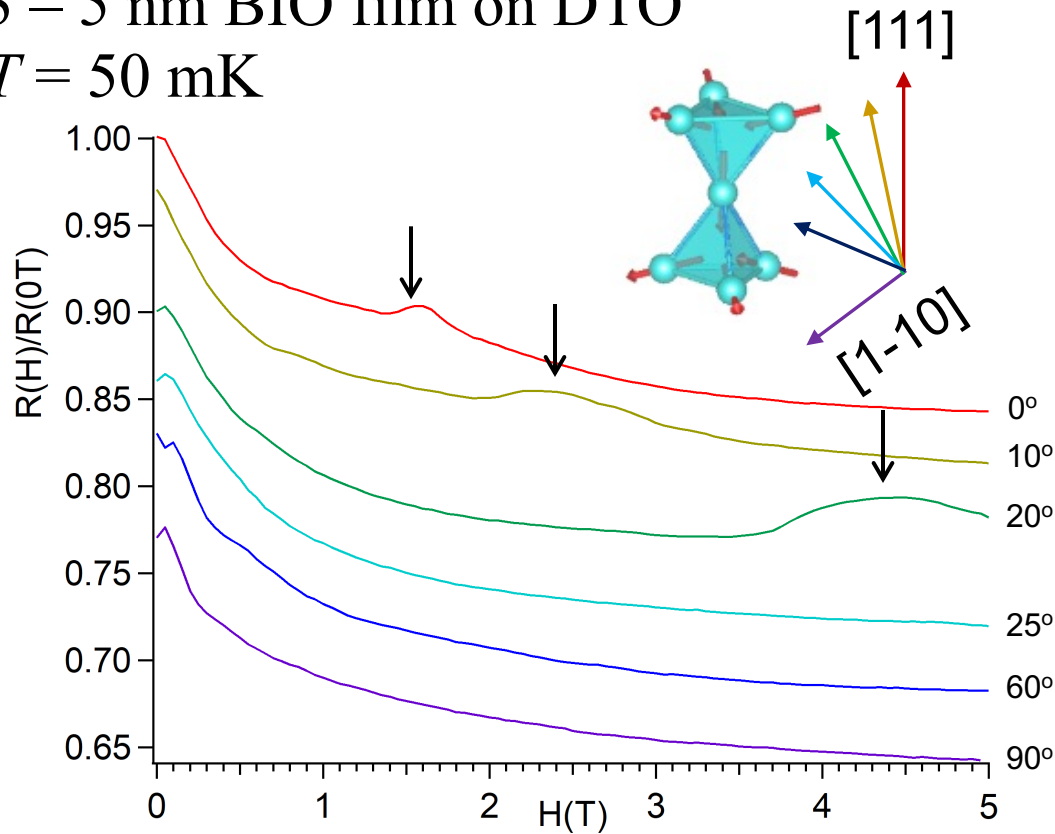


TEM, uniform interface

# Anomalous MR related to the ice-rule breaking

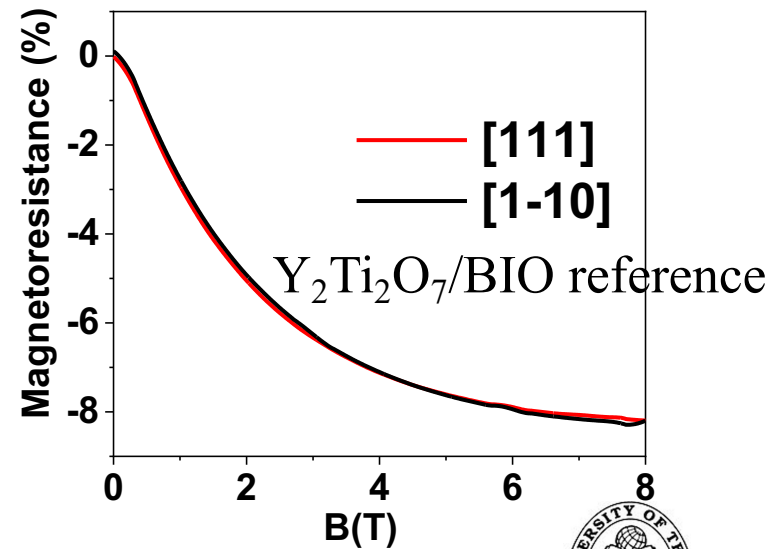
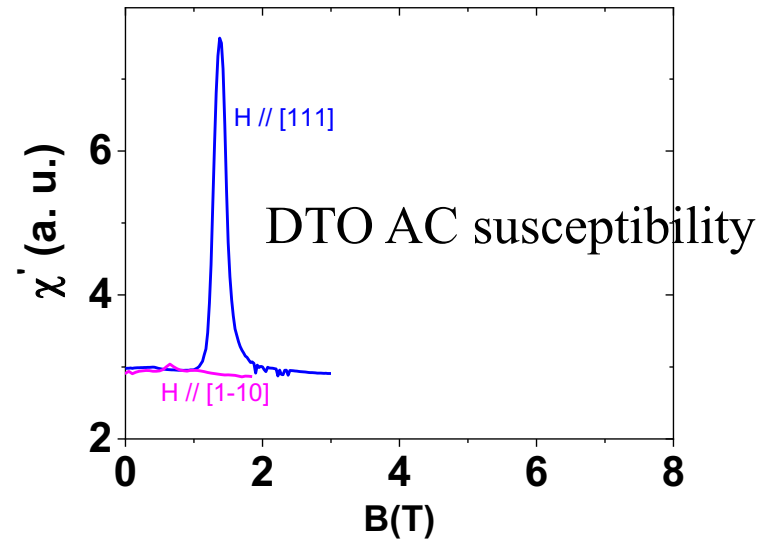
3 – 5 nm BIO film on DTO

$T = 50$  mK

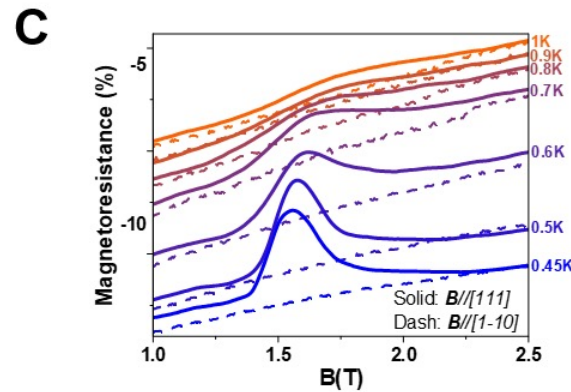
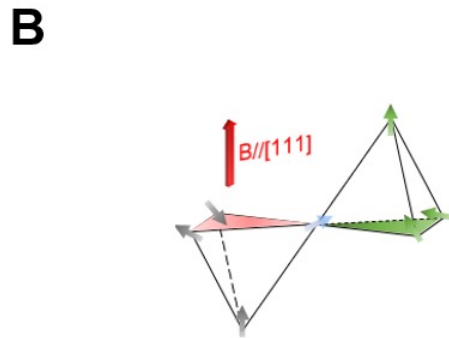
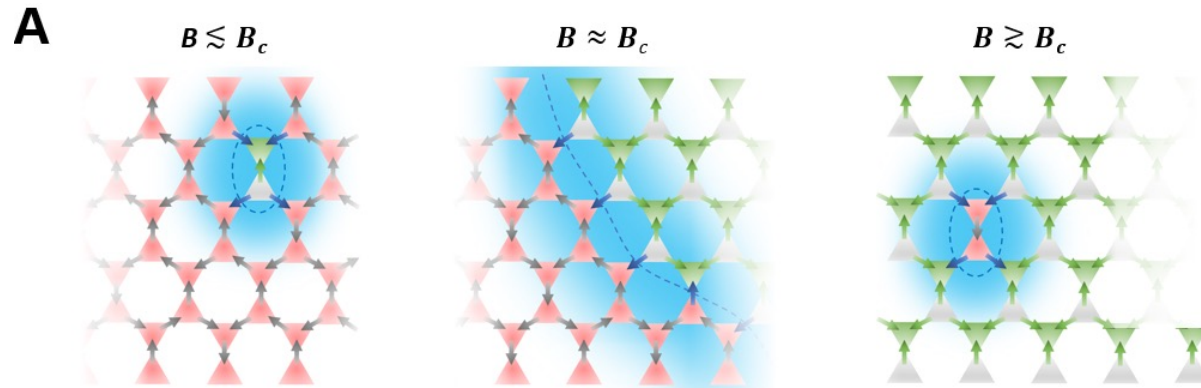


MR feature occurs while DTO enters the three in one out state

*H. Zhan, H.D. Zhou, J. Liu et al., arXiv:2011.09048*



# Anomalous MR (II): Temperature dependence

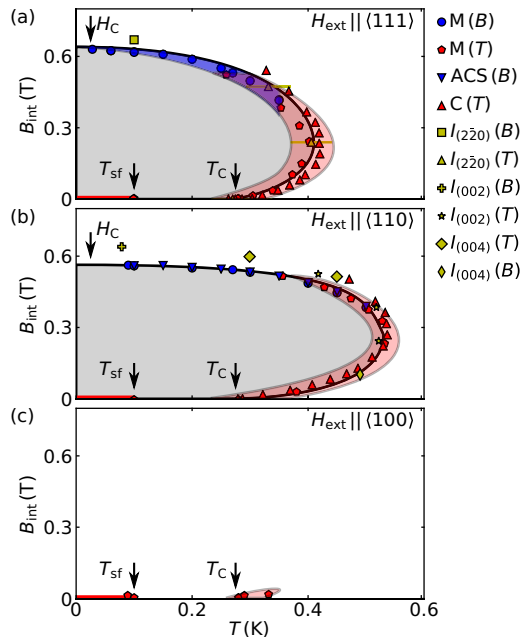
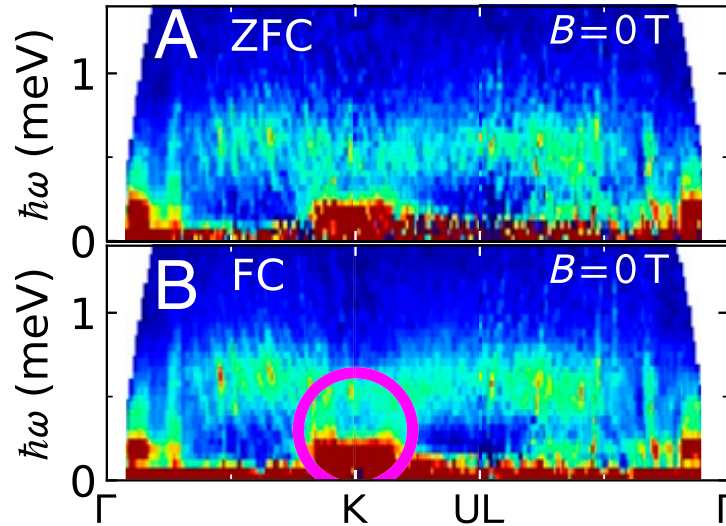
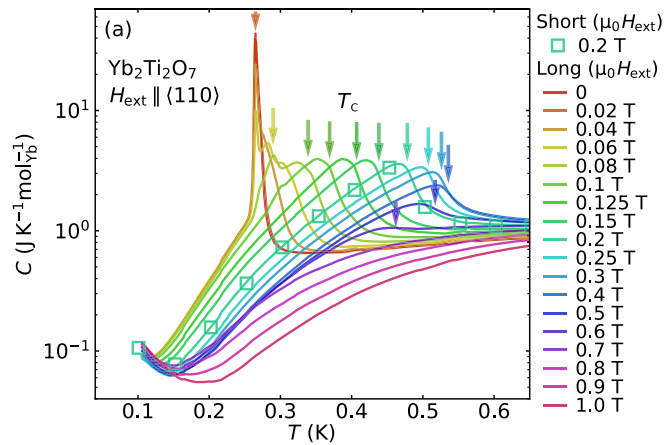


MR feature disappears above spin ice temperature

The anomalous MR responses can be depicted with the coexistence of the two spin states in the Kagome plane perpendicular to the field



# Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>/Bi<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> heterostructure



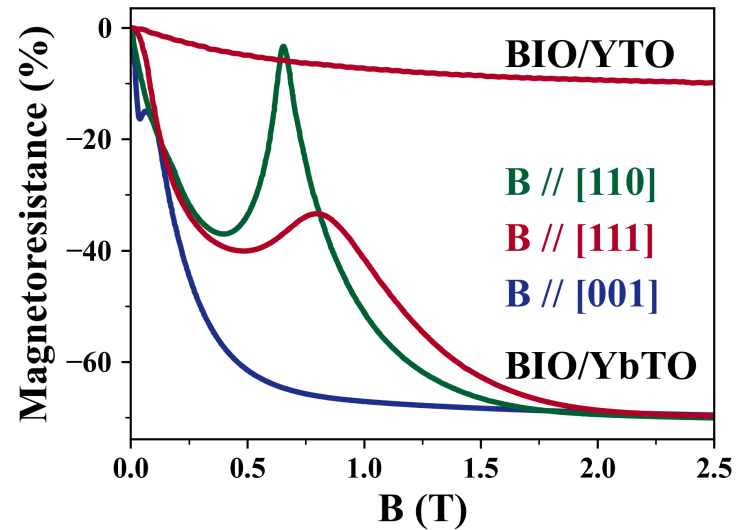
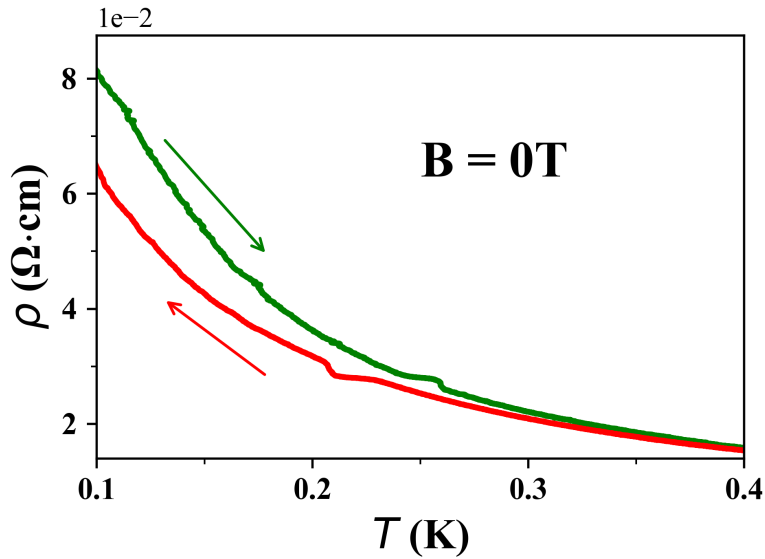
Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>: a ferromagnetic ordering at 0.28 K with strong quantum spin fluctuations

Ideal system to detect quantum spin fluctuations electronically

*A. Scheie et al., PNAS 117, 27245 (2020).*

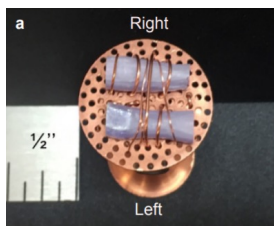
*S. Saubert et al., Physical Review B 101, 174434 (2020).*

# Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>/Bi<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> heterostructure



MR anomalies are related to strong quantum spin fluctuations at zero field, and the field dependence of the low energy excitation

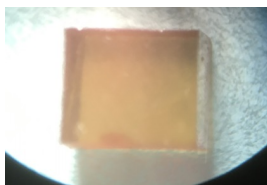
## Image Furnace



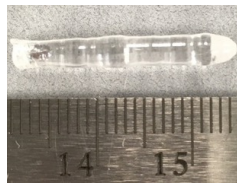
**YbMgGaO<sub>4</sub>**  
TLAF QSL



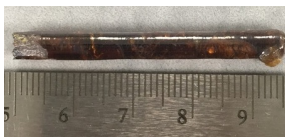
**Ba<sub>3</sub>CoSb<sub>2</sub>O<sub>9</sub>**  
TLAF QMP



**Dy<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>**  
Pyrochlore Spin ice



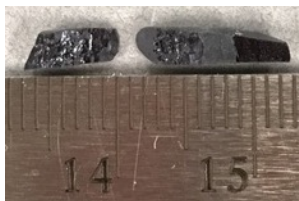
**Yb<sub>3</sub>Ga<sub>5</sub>O<sub>12</sub>**  
Garnet



**CaTiO<sub>3</sub>**  
Substrate



**SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub>**  
Magnetization Plateau



**Sr<sub>2</sub>RuO<sub>4</sub>**  
SC

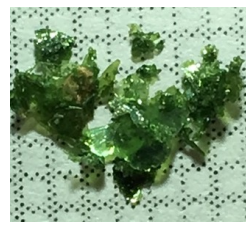


**MnO**  
Magneto-structural coupling

## Chemical Vapor Transport



**Cu<sub>2</sub>OSeO<sub>3</sub>**  
Skyrmion

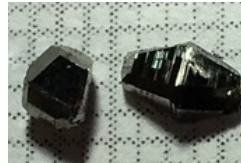


**MnPS<sub>3</sub>**  
2D magnet



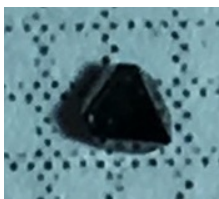
**MnP**

SC under pressure

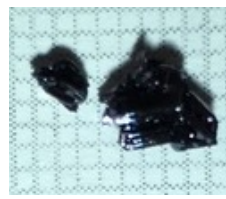


**NbP**

Weyl semimetal

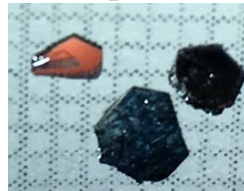


**ZnCr<sub>2</sub>Se<sub>4</sub>**  
spinel



**RuO<sub>2</sub>**

Itinerant magnet



**CoNb<sub>2</sub>S<sub>6</sub>**

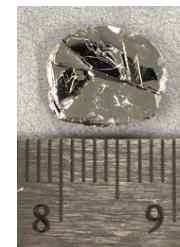
Anomalous Hall Eff.



**FeS**

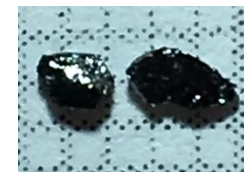
MI transition

## Flux



**NiTe<sub>2</sub>**

Topological phase? Honeycomb

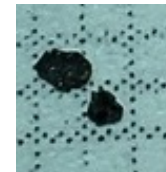


**MnB<sub>4</sub>**

MI transition?



**Er<sub>2</sub>Si<sub>2</sub>O<sub>7</sub>**



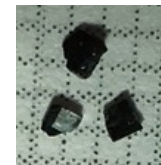
**WB<sub>2</sub>**

Hexagonal?



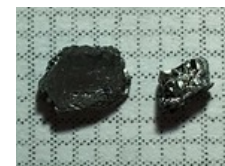
**CuCrO<sub>2</sub>**

TLAF multiferroic

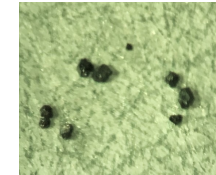


**GaFeO<sub>3</sub>**

Multiferroic?



**Co<sub>3</sub>Sn<sub>2</sub>S<sub>2</sub>**  
Magnetic Weyl



**Ba<sub>2</sub>YIrO<sub>6</sub>**  
Double Perov.

# Summary

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Geometrically frustrated lattice, magnetic ordering at extremely low temperatures, disorder, and spin-1/2, are good ingredients for QSL

Demonstrate a new route to electronically probe the exotic dynamics of geometrically frustrated quantum magnets through epitaxial interfaces.

Crystals = hard work + FedEx



*Thank You*